

the guide wavelength remains constant as the screen diameter is increased. Thus the effect of the screen on the  $HE_{11}$  mode becomes small and the wave propagates in a quasi-dielectric rod mode. However, (2) and (3) show that higher-order modes are also produced by the screen and a pure quasi-dielectric rod  $HE_{11}$  mode is not propagated. The theoretical results in Fig. 2 were taken from Beam and Wachowski.<sup>1</sup>

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### On the Measurement of Detector Impedance\*

A detector or receiver may be used in conjunction with a directional coupler and a calibrated variable short to measure its own impedance. The observed VSWR as the short is adjusted is the VSWR of the detector if the coupling is light. Various minor corrections are described.

The purpose of this paper is to describe a method for the measurement of the impedance of a detector of known response law without the use of an auxiliary detector. The method is especially suitable for use at the lowest power level within its range of operation. The detector could be a crystal, a bolometer or a receiver. The method is related to the resonance curve (Chipman method) impedance measuring technique. It is also related to the usual method of measuring impedance of a receiver in which the source and detector of a standard setup are exchanged.<sup>1</sup> The unusual feature is that the detector is used to measure its own impedance.<sup>2</sup>

It is assumed, of course, that the detector law is known and that its impedance does not change for the moderate change of received power levels encountered during the VSWR measurements. The method utilizes a directional coupler, a matched load, and a moving short with a position indicator, in addition to the signal source and the detector. A slotted line is not required. The method can be applied to standard microwave circuits but has been used primarily in

circuits utilizing plane waves in the sub-millimeter wavelength range.<sup>3</sup>

Let a signal source and a matched termination be connected to the main arms of a directional coupler, the detector to be tested to the forward-coupled side arm and a calibrated shorting (tuning) plunger to the reverse arm, as sketched in Fig. 1. Let the reflection coefficient of the mount be  $\Gamma$  and the corresponding transmission coefficient to the detector element itself be  $T$ . The voltage transmission coefficient of the coupler is  $t$  and its voltage coupling coefficient is  $r$ . The shorting plunger is assumed perfect and has a reflection coefficient  $e^{j\theta}$ .

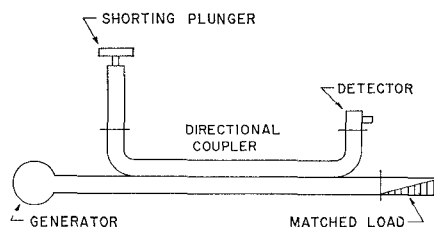


Fig. 1—Circuit for measuring detector impedance.

An incident wave of unit amplitude, upon reaching the coupler, will be partially transmitted to the detector. It will deliver a voltage  $rT$  to the detector itself. The reflection from the detector, of magnitude  $r\Gamma$ , will pass twice through the coupler on its way to the shorting plunger and back. It will contribute an additional voltage  $r\Gamma t^2 e^{j\theta} T$  at the detector. Multiple reflections will exist; the total signal delivered to the detector will be

$$E = \frac{rT}{1 - t^2 e^{j\theta} \Gamma}.$$

Maxima and minima will be observed as the shorting plunger is adjusted to make the denominator real. Their ratio is:

$$\rho = \frac{1 + t^2 |\Gamma|}{1 - t^2 |\Gamma|},$$

from which we obtain the reflection coefficient of the detector mount:

$$|\Gamma| = \frac{\rho - 1}{t^2(\rho + 1)}.$$

For light coupling,  $t \approx 1$  and  $\rho$  will be approximately the voltage standing-wave ratio of the mount.

The measurement, of course, can be no more accurate than the equipment used. In the equations above, it has been assumed that the directivity of the coupler is infinite and that the matched load is perfect. Imperfections will have an effect similar to that of the residual VSWR of a slotted line.

It will be noted that no auxiliary detector is used in this measurement. It is required only that the signal source be

moderately well-matched and that the transmission coefficient  $t$  of the coupler be known. The loss in the shorting plunger may also be lumped in with  $t$  if it has been measured. Note that the losses in the mount have not been measured. They affect the value of  $T$  which is eliminated in the ratio of maximum to minimum signals. The mount efficiency must be measured by an appropriate series of measurements of mount impedance.

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### The Synthesis of $N$ -Port Circulators\*

Professor P. Penfield has kindly brought to our notice his paper on lossless three-ports,<sup>1</sup> which independently covers some of the work described by us.<sup>2</sup> He has pointed out that for circulator synthesis from a three-port junction only one of the three conditions of equation (16) or (17) need be specified. This can be seen from the first two equations of (9) of our paper. Otherwise our classifications of three-port networks are the same.

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\* Received April 5, 1963.

<sup>1</sup> P. Penfield, Jr., "A classification of lossless three-ports," IRE TRANS. ON CIRCUIT THEORY, vol. CT-9, pp. 215-223; September, 1962.

<sup>2</sup> B. L. Humphreys and J. B. Davies, "The synthesis of  $N$ -port circulators," IRE TRANS. ON MICROWAVE THEORY AND TECHNIQUES, vol. MTT-10, pp. 551-554; November, 1962.

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<sup>1</sup> E. L. Ginzton, "Microwave Measurements," McGraw-Hill Book Co., Inc., New York, N. Y., pp. 287 ff., 307; 1957.

<sup>2</sup> The authors are indebted to R. W. Beatty for calling to their attention a related technique in which an iris mismatch at the generator and a squeezable waveguide section replace the directional coupler and the sliding short used in this paper. The technique was described by L. S. Liberman, "A method of measuring VSWR of video detectors," Radiotekhnika i elektronika, vols. 2, no. 7, pp. 941-942, 1957; Radio Engg. and Electronics, vol. 2, p. 180, 1957. (English translation.)

<sup>3</sup> R. H. Miller, P. A. Szente, and K. B. Mallory, "A Measurement of Bolometer Mount Efficiency at Millimeter Wavelengths," presented at the Millimeter and Submillimeter Conf., Orlando, Fla.; January, 1963.

### Broad-Band Microwave Discriminator\*

A sketch of a novel broad-band microwave discriminator is shown in Fig. 1. The device utilizes a pair of symmetric 3-db hybrids joined by unequal lengths of transmission line. The difference between line lengths is represented as a frequency dependent phase difference  $\phi$ . A straightforward analysis of the circuit yields

$$|E_1| = 0.707E(1 + \cos \phi)^{1/2} \quad (1)$$

$$|E_2| = 0.707E(1 - \cos \phi)^{1/2} \quad (2)$$

\* Received April 8, 1963.

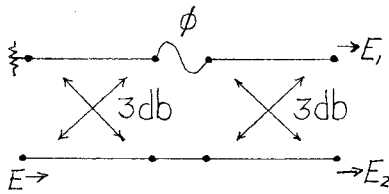


Fig. 1—Schematic of broad-band microwave discriminator.

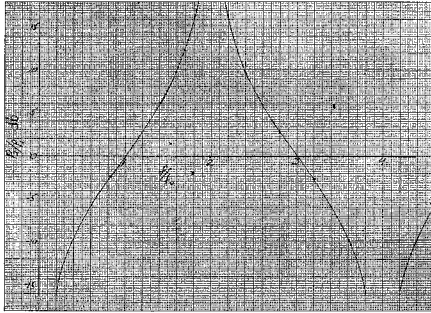


Fig. 2—Theoretical discriminator characteristic.

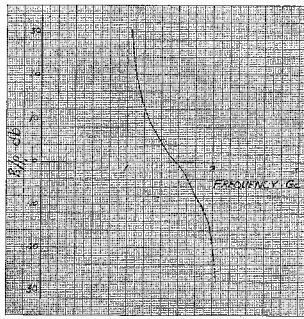


Fig. 3—Experimental data on broad-band microwave discriminator.

It is readily determined from (1) and (2) that

$$|E_1|^2 + |E_2|^2 = |E|^2 \quad (3)$$

and that the device is then 100 per cent efficient.

The ratio of output powers is

$$\frac{|E_2|^2}{|E_1|^2} = \frac{\rho_2}{\rho_1} = \tan^2 \frac{\phi}{2} \quad (4)$$

Since  $\phi$  is frequency dependent, the ratio of output power is also frequency dependent. Since  $\phi = 2\pi(l/\lambda)$ , (4) may be modified to

$$\frac{\rho_2}{\rho_1} = \tan^2 \frac{\pi l}{\lambda} \quad (5)$$

where  $l$  is the physical line length and  $\lambda$  is the wavelength in the line.

For TEM Propagation (coax), (5) may be written as

$$\frac{\rho_2}{\rho_1} = \tan^2 \pi l \frac{f}{c} \quad (6)$$

where  $c$  is the velocity of propagation in the line.

Eq. (6) may be normalized conveniently if the frequency at which  $l$  is a quarter wavelength is defined as  $f_0$ . Eq. (6) is then

$$\frac{\rho_2}{\rho_1} = \tan^2 \frac{\pi f}{4 f_0} \quad (7)$$

A plot of (7) is shown in Fig. 2.

From Fig. 2 it is plain that arbitrarily wide or narrow percentage bandwidth may be achieved by suitable choice of the line length difference. Fig. 3 shows data obtained on an experimental coaxial discriminator. The differential line length was approximately three quarter wavelengths long at 1.5 Gc.

The discriminator described is highly efficient and capable of operation over bandwidths limited only by the hybrid.

The author wishes to thank Mr. Adelsberg of the Naval Material Laboratory, Brooklyn, N. Y., for the data on the 1-2 Gc discriminator.

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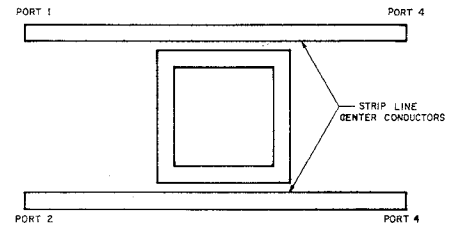


Fig. 1—Loop type strip-line directional filter.

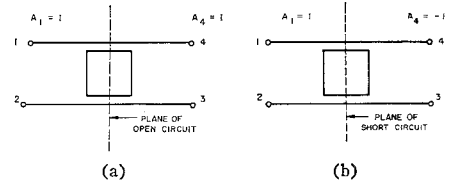
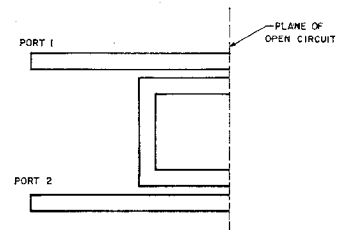
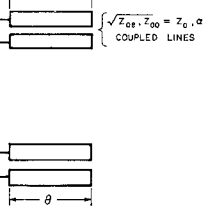


Fig. 2—Symmetrical and antisymmetrical excitation of the filter. (a) Symmetrical. (b) Antisymmetrical.

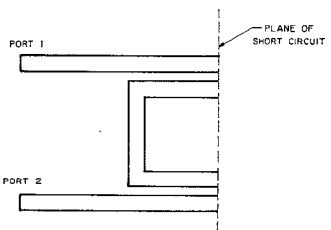


(a)

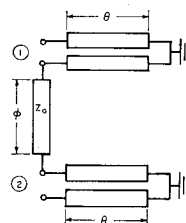


(b)

Fig. 3—Network for symmetrical excitation. (a) Two port with symmetrical excitation. (b) Equivalent circuit.



(a)



(b)

Fig. 4—Network for antisymmetrical excitation. (a) Two port with antisymmetrical excitation. (b) Equivalent circuit.

### Frequency Response of Strip-Line Traveling-Wave Directional Filters\*

Coale analyzed the single resonator strip-line traveling-wave filter shown in Fig. 1 by a perturbation method.<sup>1</sup> Design criteria were formulated relating loaded  $Q$  to pertinent circuit parameters. Experimentally it is observed that properly aligned filters exhibit a frequency response which is approximately Butterworth. The purpose of this communication is to show that the approximate frequency response of narrow bandwidth filters in the absence of dissipation and resonator discontinuities is theoretically Butterworth. Comments on the design problem are included as well as a discussion of the effects of resonator discontinuities.

The method of analysis employed makes use of the physical symmetry of the filter. By applying symmetric and antisymmetric excitation to two colinear arms of the filter as shown in Fig. 2, the four-port problem is reduced to that of solving two, two-port problems. The two, two-ports to be analyzed are shown schematically in Figs. 3 and 4. Making use of the image parameters derived by Jones and Bolljahn for the coupled lines,<sup>2</sup> and using  $ABCD$  matrix notation it is readily shown that

$$T_s = \frac{2}{2[\cos \phi \cosh 2\alpha + \sin \phi \sinh 2\alpha \cot 2\theta] + j[\sin \phi(2 \cosh^2 \alpha - \sinh^2 \alpha(\cot^2 \theta + \tan^2 \theta)) - 2 \cos \phi \sinh 2\alpha \cot 2\theta]} \quad (1)$$

$$\Gamma_s = j \frac{T_s}{2} ([\sin \phi \sinh^2 \alpha(\tan^2 \theta - \cot^2 \theta)] + \cos \phi \sinh 2\alpha(\cot \theta + \tan \theta)) \quad (2)$$

\* Received April 9, 1963.

<sup>1</sup> F. S. Coale, "A traveling-wave directional filter," IRE TRANS. ON MICROWAVE THEORY AND TECHNIQUES, vol. MTT-4, pp. 256-260; October, 1956.

<sup>2</sup> E. M. T. Jones and J. T. Bolljahn, "Coupled strip-transmission-line filters and directional couplers," IRE TRANS. ON MICROWAVE THEORY AND TECHNIQUES, vol. MTT-4, pp. 75-81; April, 1956.